# A Hierarchical Geometric Framework to Design Locomotive Gaits for Highly Articulated Robots 

Author Names Omitted for Anonymous Review. Paper-ID [204]


#### Abstract

Motion planning for mobile robots with many degrees-of-freedom ( $\mathbf{D o F}$ ) is challenging due to their highdimensional configuration spaces. To manage this curse of dimensionality, this paper proposes a new hierarchical framework that decomposes the system into sub-systems (based on shared capabilities of DoFs), for which we can design and coordinate motions. Instead of constructing a high-dimensional configuration space, we establish a hierarchy of two-dimensional spaces on which we can visually design gaits using geometric mechanics tools. We then coordinate motions among the two-dimensional spaces in a pairwise fashion to obtain desired robot locomotion. Further geometric analysis of the two-dimensional spaces allows us to visualize the contribution of each sub-system to the locomotion, as well as the contribution of the coordination among the sub-systems. We demonstrate our approach by designing gaits for quadrupedal robots with different morphologies, and experimentally validate our findings on a robot with a long actuated back and intermediate-sized legs.


## I. Introduction

High degree-of-freedom (HDoF) mobile robots offer advantages over conventional wheeled vehicles in difficult-tonegotiate terrains [3, 17, 21, 25]. However, the benefit of many degrees of freedom also poses a challenge: how to coordinate all of the DoFs to produce purposeful motion. One approach to address this so-called curse of dimensionality is to couple groups of the DoFs into a one or two DOFs thereby producing a lower dimensional configuration space [5, 10, 15, 16, 24, 2, 6]. We term lower dimensional systems as sub-systems. Sub-systems can be formed based on observations of biological systems, empirical experiments, or intuition. The approach described in this paper first plans for motions within every sub-system, and then for coordination among the sub-systems. Specifically, we propose a hierarchical framework to decompose the motion planning problem in the full configuration space of a robot, into i) a series of simpler motion planning problems in the low-dimensional spaces associated with the different sub-systems, and ii) the coordination of these sub-systems.

Our framework relies on a motion planning approach, geometric mechanics (GM), which offers powerful tools to design, visualize, and analyze a wide range of locomotor behaviors [4, 7, 14]. We rely on GM tools to visually design the motion of each subsystem with respect to our locomotive goals (here, to maximize stride displacement), and to design the coordination of pairs of sub-systems. By relying on a coordinate-descentlike approach, we then optimize locomotion by iterating over every pair of sub-systems until convergence. Once this process has converged, we further show how GM tools allows us to


Fig. 1. Three link swimmer with four cubic legs CAD diagram of the robot (a) top and (b) side view. The upper back and lower back are actuated by two AX12 servos (labeled as 1 and 2) and the legs are actuated by XL320 servos. The legs have rectangular shape $(24 \times 24 \times 45 \mathrm{~mm})$ and the body is always in contact with ground via blocks ( $55 \times 12 \times 45 \mathrm{~mm}$ ) attached to the body segments. (c) Physical model.
visualize the contribution of each sub-system to locomotion, as well as that of the coordination of every pair of sub-systems.

We applied our framework to study the motion of a quadrupedal robot with a long articulated body and intermediate-sized legs (Fig. 1), a morphology incorporating features from both lizards and snakes. There, we show how our framework provides insight into the role of leg contacts and body undulation in quadruped locomotion. Previous work has shown that the undulation of an elongated body improves the stride length of animals and robots [11, 27, 8]. There are two commons modes of body undulation: standing waves (lizardlike) [11, 27], and traveling waves propagated from head to tail (snake-like) [8]. In the first mode, the body curvature is maximized to increase the self-propulsion generated by legground interaction. In the other mode, the body undulation can generate self-propulsion even without legs. Interestingly, we find that the optimized whole-body motion (i.e., motion in the full configuration space) for our robot is a hybrid combination of the lizard- and snake-like modes, such that the self-propulsion from leg-ground contact and that from body-ground contact are properly coordinated. We validate our theoretical predictions through a series of experiments where a $10-$ DoF robot locomotes on granular material.
The paper is structured as follows: Section $\Pi$ introduces geometric mechanics and presents our framework to design whole-body motion; Section III applies our methods to study the motion design of a robot with a long articulated body and four legs; Section IV discusses the significance and conclusions of our work.


Fig. 2. Example height functions and gaits in different shape spaces (a) The height function in a Euclidean shape space corresponds to motions of an 8-link snake robot slithering in the forward direction on the surface of 6 mm plastic particles. The purple circle represents the optimal gait in the corresponding shape space. (b) The height function in a torus shape space corresponds to the forward motion of a robot with a long actuated body and intermediate-sized leg on the surface of $\sim 1 \mathrm{~mm}$ poppy seeds. The solid purple curve represents a sample gait in the corresponding shape space. Orange lines represent the assistive lines to form closed loops with the gait path in the unfolded torus shape space. The area in the solid purple shadow represents the area where the gait path and the assistive lines form a clockwise loop; the area in the dashed purple shadow represents the area where the gait path and the assistive lines form a counterclockwise loop. Red, white and black indicate positive, zero and negative values in height function respectively.

## II. Method

## A. Review of geometric mechanics

In this section, we provide a concise overview of the geometric tools needed for this paper. For a more detailed and comprehensive review, we refer readers to [4, 7, 14]. The gait design tools of geometric mechanics separates the configuration space of a system into two spaces: the position space and the shape space. The position space denotes the location of a system relative to a world frame, and the shape space denotes the internal configuration (internal shape) of the system. Geometric mechanics seeks to establish a functional relationship between velocities in these spaces; this functional relationship is often called a connection, and it shares many properties with the Jacobian of a robotic manipulator.

1) Kinematic reconstruction equation: Principally, in kinematic systems where there is no drift in the system, the equations of motion [14] reduce to

$$
\begin{equation*}
\xi=A(r) \dot{r}, \tag{1}
\end{equation*}
$$

where $\boldsymbol{\xi}=\left[\begin{array}{lll}\xi_{x} & \xi_{y} & \xi_{\theta}\end{array}\right]^{T} \in g$ denotes the body velocity in forward, lateral and rotational direction; $\boldsymbol{r}$ denotes the internal shape variables (joint angles); $\boldsymbol{A}(\boldsymbol{r})$ is the local connection matrix that relates shape velocity $\dot{\boldsymbol{r}}$ to body velocity $\boldsymbol{\xi}$. Eq. (1) is also called the kinematic reconstruction equation and maps the changes in internal shape variables (joint angles) to changes in group variables (position and orientation) of the robot.

The interaction between a robot and an environment is often difficult to model, which is typically the case when the system locomotes on granular medium. In such a case, calculating the local connection matrix $\boldsymbol{A}$ can be quite challenging. Instead, we numerically derive $\boldsymbol{A}$ using resistive force theory
(RFT) [13, 22, 26] to model the contact. Prior work [8] has shown that numerically derived local connections using granular RFT can effectively predict movements in granular media. In this paper, we model the robot-ground contact by poppy seed RFT equations [27], from which we numerically derive the local connections.
2) Connection vector fields and height functions: Each row of local connection matrix $\boldsymbol{A}$ corresponds to a component direction of the body velocity and therefore gives rise to a connection vector field. The body velocities in the forward, lateral and rotational directions are respectively computed as the dot product of connection vector fields and the shape velocity $\dot{\boldsymbol{r}}$. A shape velocity $\dot{\boldsymbol{r}}$ along the direction of the vector field yields the largest possible body velocity, while a shape velocity $\dot{\boldsymbol{r}}$ orthogonal to the field produces zero body velocity.
A gait can be represented as a closed curve in the corresponding shape space. The displacement resulting from a gait, $\partial \chi$, can be approximated by

$$
\left(\begin{array}{l}
\Delta x  \tag{2}\\
\Delta y \\
\Delta \theta
\end{array}\right)=\int_{\partial \chi} A(r) d r
$$

According to Stokes' Theorem, the line integral along a closed curve $\partial \chi$ is equal to the surface integral of the curl of $\boldsymbol{A}(\boldsymbol{r})$ over the area enclosed by $\partial \chi$ :

$$
\begin{equation*}
\int_{\partial \chi} A(r) d r=\iint_{\chi} \nabla \times A(r) d r_{1} d r_{2} \tag{3}
\end{equation*}
$$

where $\chi$ denotes the area enclosed by $\partial \chi$. The curl of the connection vector field, $\boldsymbol{\nabla} \times \boldsymbol{A}(\boldsymbol{r})$, is referred to as the height function (Fig. 2 a ). The three rows of the vector field $\boldsymbol{A}(\boldsymbol{r})$ can thus produce three height functions in the forward, lateral and rotational direction, respectively.
With the above derivation, we simplify the gait design problem to drawing a closed path in Euclidean shape space. The displacements can be approximated by the volume enclosed by the gait path. For example, in Fig. 2 a , we show an example height function in Euclidean shape space corresponding to the motion in the forward direction of an 8 -link snake robot slithering on the surface of 6 mm plastic particles. The circle (solid purple circle in Fig. 2a) enclosing the most volume represents the gait with the largest forward stride displacement.

Although the review of geometric mechanics section is limited in two dimensions, there has been recent work on applying geometric mechanics to high-dimensional robot gait design [5, 20]. Ramasamy and Hatton [20] extended geometric mechanics to high dimensional shape space and proposed a method to visualized the kinematics. Gong et al. [5] approximated the kinematics of complex locomotors using only two shape variables with the proper shape basis functions.
3) Torus shape space: Often, a shape space can have some cyclic structure. The shape space with both shape variables being cyclic, $\left(T^{2}\right)$, has a torus shape [12]. In Fig. 2p, we show an example height function in a torus shape space
corresponding to the forward motion of a quadruped robot with the standing wave body undulation on the surface of $\sim 1$ mm poppy seeds. The shape variable, $\boldsymbol{\Phi}=\left[\phi_{1}, \phi_{2}\right]^{T} \in T^{2}$, has the cyclic structure, where one axis represents the phase of the leg movement, and the other axis represents the phase of the body undulation. While the gait path in (solid purple curve Fig. 2 p ) is a closed curve in the torus shape space, there is no obvious area enclosed by the gait path.

To form an enclosed area, we introduce the notion of two assistive lines in the Euclidean parameterization of the torus shape space. The assistive lines are defined as: $\partial \chi_{1}: \phi_{1}=0$, (the solid orange line $l_{1}+l_{2}$ in Fig. 2 p ), and $\partial \chi_{2}: \phi_{2}=0$, (the solid orange line $l_{3}+l_{4}$ in Fig. 2 p ). Note that the dashed line $l_{i}^{\prime}$ is equivalent to the corresponding solid line $l_{i}$, according to the cyclic property of the torus shape space. In Fig. 2 p , the gait path (the solid purple curve) together with two assistive lines $\left(l_{1}^{\prime}, l_{2}, l_{3}^{\prime}, l_{4}\right)$ encloses two areas in the torus shape space (solid purple shadow area and dashed purple shadow area). The net displacement can be approximated by the line integral along two assistive lines plus the surface integral of the enclosed area:

$$
\begin{array}{r}
\int_{\partial \chi} A(r) d \Phi=\int_{\partial \chi_{1}} A(\Phi) d \Phi+\int_{\partial \chi_{2}} A(\Phi) d \Phi+ \\
\iint_{\chi-\chi_{1}-\chi_{2}} \nabla \times A(\Phi) d \phi_{1} d \phi_{2} \tag{4}
\end{array}
$$

where $\chi-\chi_{1}-\chi_{2}$ denotes the area enclosed by the assitive lines and gait path. Note that in Fig. 2 p , the gait path and the assistive lines form a clockwise loop in the lower right corner (solid purple shadow area) and a counterclockwise loop in the upper left corner (dashed purple shadow area). Taking the handedness into account, the surface integral (third term on RHS) in Eq. (4) should be computed as the surface integral of enclosed area in the clockwise loop (solid purple shadow area in Fig. 2p) minus the surface integral of enclosed area in the counterclockwise loop (dashed purple shadow area in Fig. 2 b ).

To interpret the physical meaning of each term (RHS) in Eq.(4), we introduce the notation of two sets $P_{a}$ and $P_{b}$, where $P_{a}$ and $P_{b}$ denotes the contribution from the horizontal and vertical coordinate axis in the torus shape space. In our example in Fig. $2 \mathrm{p}, P_{a}$ represents the contribution from leg movement and $P_{b}$ represents the contribution from standing wave body undulation.

The physical meaning of the first term (RHS) in Eq. (4) is the displacement solely resulting from motions in $\phi_{2}$, while keeping $\phi_{1}$ constant at initial position, i.e., $P_{a}^{C} \cap P_{b}$, where ()$^{C}$ denotes the complement set. The physical meaning of the second term (RHS) in Eq. (4) is the displacement solely resulting from motions in $\phi_{1}$, while keeping $\phi_{2}$ constant at initial position, i.e. , $P_{a} \cap P_{b}^{C}$. The first two terms (RHS) in Eq.(4) are constant and independent from any gait path we choose. The physical meaning of the last term (RHS) in Eq. (4) is the additional displacement resulting from the contribution from coordination of $\phi_{1}$ and $\phi_{2}$, i.e., $P_{a} \cap P_{b}$.

Finally, we refer reader to [7] for detailed derivation and proof of motion planning in torus shape spaces.

## B. Whole-body motion

1) Sub-system motion phase variable: In this section, we construct a high-dimensional torus shape space, on which we can design the whole-body coordination. Consider an $M$-dimensional motion system and its shape variable $r=$ $\left[r_{i}, r_{2}, \ldots, r_{M}\right]^{T} \in R^{M}$. Assume there are $N$ sub-systems. We term the motion of a sub-system sub-system motion. Each sub-system motion will form a closed loop gait path in the shape space $R^{M}$. The closed loop gait path in sub-system motion $i \in\{1,2, \ldots, N\}$ can be described by a function $\boldsymbol{f}_{i}: T^{1} \rightarrow R^{M}$ such that $\boldsymbol{r}=\boldsymbol{f}_{\boldsymbol{i}}\left(\phi_{i}\right)$, where $\phi_{i} \in T^{1}$ is the gait phase variable of sub-system motion $i$. The wholebody motion can be formulated as the superposition of the sub-system motions gait paths:

$$
\begin{equation*}
\boldsymbol{r}=\sum_{i=1}^{i=N} \boldsymbol{f}_{\boldsymbol{i}}\left(\phi_{i}\right) \tag{5}
\end{equation*}
$$

We define the sub-system motion phase variable as $\boldsymbol{\Phi}=$ $\left[\phi_{1}, \phi_{2}, \ldots, \phi_{N}\right]^{T} \in T^{N}$. We can relate the $\dot{\Phi}$ to the shape velocity $\dot{\boldsymbol{r}}$ by the Jacobian matrix [23] $\boldsymbol{J}_{r, \Phi}$ :

$$
\dot{\boldsymbol{r}}=\left[\begin{array}{llll}
\frac{\boldsymbol{d} \boldsymbol{f}_{1}\left(\phi_{1}\right)}{\boldsymbol{d} \phi_{1}} & \frac{\boldsymbol{d} \boldsymbol{f}_{2}\left(\phi_{2}\right)}{\boldsymbol{d} \phi_{2}} & \ldots & \frac{\boldsymbol{d} \boldsymbol{f}_{N}\left(\phi_{N}\right)}{\boldsymbol{d} \phi_{N}} \tag{6}
\end{array}\right] \dot{\boldsymbol{\Phi}}=\boldsymbol{J}_{\boldsymbol{r}, \boldsymbol{\Phi}} \dot{\boldsymbol{\Phi}}
$$

With Eq. (6) in hand, we can rewrite Eq.(1):

$$
\begin{align*}
& \xi=A(r) \dot{r} \\
& =A\left(\sum_{i=1}^{i=N} f_{i}\left(\phi_{i}\right)\right) \dot{r}=A\left(\sum_{i=1}^{i=N} f_{i}\left(\phi_{i}\right)\right) J_{r \Phi} \dot{\Phi} \\
& =A^{\prime}(\Phi) \dot{\Phi} \tag{7}
\end{align*}
$$

where $A^{\prime}(\Phi)$ is the new local connection matrix mapping the differential of sub-system motion phase variable, $\dot{\mathbf{\Phi}}$, to the body velocity $\boldsymbol{\xi}$.

A properly designed gait in the sub-system motion phase variable space ( $N$-dimensional torus shape space) will then lead to coordinated whole-body motion. That is, we simplify the whole-body motion gait design problem in the original $M$ dimensional Euclidean space $R^{M}$ to the $N$-dimensional torus sub-system motion phase variable space $T^{N}$.
2) Gait design: In this section, we will propose a method to design a gait path in an $N$-dimensional torus shape space. The gait path in the $N$-dimensional sub-system motion phase variable space can be described by a gait function $H: T^{1} \rightarrow$ $T^{N}$, such that: $\boldsymbol{\Phi}=\left[h_{1}\left(\phi_{0}\right), h_{2}\left(\phi_{0}\right), \ldots, h_{N}\left(\phi_{0}\right)\right]^{T} . \phi_{0} \in T^{1}$ is the global gait phase of the motion in sub-system motion phase variable space.

We assume that all the functions $h_{i}$ are homeomorphism mapping, such that $h^{-1} \circ h$ is the identity mapping. Under this assumption, all sub-system motions will operate at the same frequency (more formally, the winding number of the gait path must be $[1,1, \ldots, 1]$ in $N$-dimensional torus space).

While there are $N$ functions in $H$, there are only $N-1$ independent functions in $H$ (dimension of $\Phi \in T^{N}$ space minus the dimension of $\phi_{0} \in T^{1}$ space), allowing one constraint equation on $H$. For simplicity, we fix $h_{1}(x):=x$, i.e., $\phi_{1}:=\phi_{0}$. Note that our method is permutation invariant.

For simplicity, in this work we assume that all the functions in $H$ are linear mapping, such that $h_{i}(x)=x+\phi_{i}$, where $\phi_{i}$ is the phase offset of the sub-system motion $i$. In these cases, only the phase offsets between sub-system motions are subject to optimization in the whole-body coordination gait design. For more complex systems where linear mapping cannot produce the effective whole-body coordination, the mapping function prescription can be extended to sigmoidal, quadrutic, cubic or other homeomorphism mappings.

We start the optimization with a naive initial guess on the gait function $H$ that all functions in $H$ are identity mapping, i.e. , $h_{i}(x)=x$. We then decompose the $N$-dimensional torus shape space into $N-1$ two-dimensional torus shape spaces $\left[\phi_{0}, \phi_{i}\right]^{T}$, where $i \in\{2,3, \ldots N\}$. Next, we optimize one independent function $h_{i} \subset H$ in the two-dimensional torus shape space composed of $\left[\phi_{0}, \phi_{i}\right]^{T} \in T^{2}$. Then, we take an approach similar to coordinate descent methods, where at each step, we optimize and update one function in $H$ and iterate among all the independent functions in $H$ until convergence. In the converged gait function $H^{*}$, every independent function $h_{i}^{*} \subset H^{*}$ should be optimized in its corresponding twodimensional torus shape space composed of $\left[\phi_{0}, \phi_{i}\right]^{T}$.

The most challenging part in our proposed gait optimization tool is to find the optimized mapping $h_{i}$ in the two-dimensional torus shape space $\left[\phi_{0}, \phi_{i}\right]^{T}$. The detailed optimization steps are presented in the remainder of this section.

To optimize each independent function $h_{i} \subset H$, we first relate the two-dimensional shape variable $\left[\phi_{0}, \phi_{i}\right]^{T}$ to the sub-system motion phase variable $\boldsymbol{\Phi}$ such that $\boldsymbol{\Phi}=$ $\left[h_{1}\left(\phi_{0}\right), h_{2}\left(\phi_{0}\right), \ldots, h_{i-1}\left(\phi_{0}\right), \phi_{i}, h_{i+1}\left(\phi_{0}\right), \ldots, h_{N}\left(\phi_{0}\right)\right]^{T}$. Next, we relate the differentials of $\phi_{0}$ and $\phi_{i}$ to the subsystem motion phase variable differential by a Jacobian matrix $\boldsymbol{J}_{\boldsymbol{\Phi},\left[\phi_{0}, \phi_{i}\right]^{T}}:$

$$
\left.\dot{\boldsymbol{\Phi}}=\boldsymbol{J}_{\boldsymbol{\Phi},\left[\phi_{0},\right.} \phi_{i}\right]^{T}\left[\begin{array}{c}
\dot{\phi}_{0}  \tag{8}\\
\dot{\phi}_{i}
\end{array}\right]
$$

With Eq. (8) in hand, we can rewrite Eq. (7) as

$$
\begin{align*}
&\left.\boldsymbol{\xi}=\boldsymbol{A}^{\prime}(\boldsymbol{\Phi}) \dot{\boldsymbol{\Phi}}=\boldsymbol{A}^{\prime}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{T}\right) \boldsymbol{J}_{\boldsymbol{\Phi},\left[\phi_{0},\right.} \phi_{i}\right]^{T}\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
&=\boldsymbol{A}_{\boldsymbol{H}}^{[i]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{T}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \tag{9}
\end{align*}
$$

where $\boldsymbol{A}_{\boldsymbol{H}}^{[i]}$ is the new local connection matrix. The subscript $H$ denotes that we need the gait function $H$ to derive the local connection matrix. The superscript $[i]$ denotes we separated $h_{i}$ from $H$ in deriving the local connection matrix.

Following the method discussed in Section II-A, we construct the height function (given by $\boldsymbol{\nabla} \times \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}$ ), with the help of which we finally design the mapping $h_{i}$.
3) Visualization: In this subsection, we develop the tools to visualize the contribution from a specific sub-system motion and the contribution from the coordination of two sub-system
motions. We define $P_{i}$ as the set denoting the contribution from a sub-system motion $i$. The contribution from coordination between sub-system motion $i$ and $j$ can be expressed as $P_{i} \cap P_{j}$.

In Section II-B2, we construct the local connection matrix mapping the phase differentials $\dot{\phi}_{0}$ and $\dot{\phi}_{i}$ to the body velocity $\xi$. Following the approach introduced in Section II-A we can approximate the net displacement as:

$$
\begin{align*}
& \int_{\partial \chi} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}\left(\left[\phi_{0} \phi_{i}\right]^{T}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& \quad=\int_{\partial \chi_{1}: \phi_{0}=0} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& \quad+\int_{\partial \chi_{2}: \phi_{i}=0} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& \quad+\iint_{\chi-\chi_{1}-\chi_{2}} \nabla \times \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}\left(\left[\phi_{0} \phi_{i}\right]^{\boldsymbol{T}}\right) \boldsymbol{d} \phi_{0} \boldsymbol{d} \phi_{i} \tag{10}
\end{align*}
$$

The physical meaning of each term (RHS) Eq. 10p can be interpreted in a similar manner to our analysis in Section II-A The physical meaning of the first term (RHS) in Eq. (10) is the displacement solely from the motion in sub-system $i$ while keeping $\phi_{0}$ constant at initial position. Note that all other sub-system motions have been coupled to $\phi_{0}$ by $H^{*}$ in our derivation of local connection matrix $\boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}$. Keeping $\phi_{0}$ constant at initial position is equivalent to keeping all other sub-system motions stopped at initial position. Thus, the first term (RHS) in Eq. (10) can then be interpreted as the displacement solely from sub-system motion $i$, while keeping all other motion systems stopped at initial position, i.e. , $P_{i} \cap P_{[i]}^{C}$, where $P_{[i]}=P_{1} \cup P_{2} \cup \ldots \cup P_{i-1} \cup P_{i+1} \cup \ldots \cup P_{N}$. The subscript [i] in $P_{[i]}$ denotes that $P_{[i]}$ is the union of the contribution of all other sub-system motion except sub-system motion $i$. Since only one sub-system motion is involved in the first term (RHS) of Eq. (10), it is then a pre-computed constant, independent from the sub-system motion coordination.

The physical meaning of the third term (RHS) in Eq. (10) is the additional displacement from coordination between subsystem motion $i$ and all other sub-system motions, i.e., $P_{i} \cap$ $P_{[i]}$.

Since $P_{i}=\left(P_{i} \cap P_{[i]}^{C}\right) \cup\left(P_{i} \cap P_{[i]}\right), P_{i}$ can be visually analyzed by the surface integral over the height function $\boldsymbol{\nabla} \times \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}$ plus a pre-computed constant.

Next, we develop the tool to visualize the contribution from coordination of two specific sub-system motions $i$ and $j$, i.e., $C_{i} \cap C_{j}$. In order to construct $C_{i} \cap C_{j}$, we first assign $h_{j}=0$ such that in this scenario, sub-system motion $j$ is stopped. Following our approach in section II-B2, we construct a twodimensional shape variable $\left[\phi_{0}, \phi_{i}\right]$ and relate it to the subsystem motion phase variable $\boldsymbol{\Phi}^{[j]}$ (the superscript $[j]$ denotes that $\phi_{j}$ is not involved in $\boldsymbol{\Phi}^{[j]}$ ) as:

$$
\begin{align*}
\boldsymbol{\Phi}^{[j]}=\left[h_{1}^{*}\left(\phi_{0}\right),\right. & \ldots, h_{i-1}^{*}\left(\phi_{0}\right), \phi_{i}, h_{i+1}^{*}\left(\phi_{0}\right), \ldots, \\
& \ldots,  \tag{11}\\
h_{j-1}^{*}\left(\phi_{0}\right), & 0, \\
h_{j+1}^{*}\left(\phi_{0}\right), & \left.h_{N}^{*}\left(\phi_{0}\right)\right]^{T}
\end{align*}
$$

Given Eq.(11), we can relate the differential of $\phi_{0}$ and $\phi_{i}$ to the sub-system motion phase variable differential $\Phi^{[j]}$ by a Jacobian matrix $\boldsymbol{J}_{\boldsymbol{\Phi}}{ }^{[j]},\left[\phi_{0}, \phi_{i}\right]^{T}$ :

$$
\dot{\Phi}^{[j]}=J_{\Phi^{[j]},\left[\phi_{0}, \phi_{i}\right]^{T}}\left[\begin{array}{c}
\dot{\phi}_{0}  \tag{12}\\
\dot{\phi}_{i}
\end{array}\right]
$$

Note that motion associated with the sub-system motion phase variable $\Phi^{[j]}$ in Eq.(12) does not include the motion in sub-system $j$. Similar to Eq.(7), we can then relate it to the shape velocity:

$$
\begin{align*}
& \boldsymbol{\xi}^{[j]}=\boldsymbol{A}\left(\sum_{k=1}^{k=j-1} f_{k}\left(\phi_{k}\right)+\right.\left.\sum_{k=j+1}^{k=N} f_{k}\left(\phi_{k}\right)\right) \dot{r} \\
&=\boldsymbol{A}^{\prime}\left(\boldsymbol{\Phi}^{[j]}\right) \boldsymbol{\Phi}^{[j]}=\boldsymbol{A}^{\prime}\left(\left[\phi_{0} \phi_{i}\right]^{T}\right) \boldsymbol{J}_{\boldsymbol{\Phi}^{[j]},\left[\phi_{0}, \phi_{i}\right]^{T}}\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
&=\boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\phi_{0} \phi_{i}\right]^{T}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \tag{13}
\end{align*}
$$

where $\boldsymbol{\xi}^{[j]}$ denotes the body velocity (the superscript $[j]$ in $\xi^{[j]}$ denotes that sub-system motion $j$ is not coupled in the body velocity) and $A_{H^{*}}^{[i[j]}$ denotes the new local connection. The superscript $[i][j]$ in $A_{H^{*}}^{[i] j]}$ denotes that we separated $h_{i}$ and $h_{j}$ from $H^{*}$ in deriving $A_{H^{*}}^{[i][j]}$. With Eq.(13), we can rewrite the Eq.(3) as:

$$
\begin{align*}
& \int_{\partial \chi} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& =\int_{\partial \chi_{1}: \phi_{0}=0} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& \quad+\int_{\partial \chi_{2}: \phi_{i}=0} \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right)\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right] \\
& \quad+\iint_{\chi-\chi_{1}-\chi_{2}} \nabla \times \boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\begin{array}{ll}
\phi_{0} & \phi_{i}
\end{array}\right]^{\boldsymbol{T}}\right) \boldsymbol{d} \phi_{0} \boldsymbol{d} \phi_{i} \tag{14}
\end{align*}
$$

Similar to our analysis in Section II-B3, we can interpret the physical meaning of the third term (RHS) in Eq.(14) as the additional displacement resulting from the coordination between sub-system motion $i$ and all other sub-system motions coupled to $\phi_{0}$ in Eq.(11), where we assume that sub-system motion $j$ is stopped. In this way, we can express the physical meaning of the third term (RHS) in Eq. (14] as $P_{i} \cap\left(P_{[i]}-P_{j}\right)$.

Given Eq. (10) and Eq. (14), we can construct $P_{i} \cap P_{j}$. We subtract the third term (RHS) in Eq. (14) from the third term (RHS) in Eq. 10), we will have:


Fig. 3. The model for our robot and its sub-system motions. (a) The shape variables of a robot with a long articulated body and four legs. The body undulation joint angle, $\alpha_{i}$, and leg shoulder joint angle $\beta_{i}$ are labelled. Segments in solid ellipse are in contact with ground; segments in open ellipse are in air. (b) Leg sub-system motion. The leg sub-system motion is prescribed by a periodic variable $\phi_{1}$. The row of eight boxes indicate the contact state of the leg during the phase of the gait, where filled gray means contact and open white means in the air. The blue curves indicate the leg shoulder joint angles. (c) Upper back sub-system motion. The upper back sub-system motion is prescribed by a periodic variable $\phi_{2}$. (d) Lower back sub-system motion. The lower back sub-system motion is prescribed by a periodic variable $\phi_{3}$

$$
\begin{aligned}
& P_{i} \cap P_{j}=P_{i} \cap P_{[i]}-P_{i} \cap\left(P_{[i]}-P_{j}\right) \\
= & \iint_{\chi-\chi_{1}-\chi_{2}} \nabla \times\left(\boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i]}\left(\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right]\right)-\boldsymbol{A}_{\boldsymbol{H}^{*}}^{[i][j]}\left(\left[\begin{array}{c}
\dot{\phi}_{0} \\
\dot{\phi}_{i}
\end{array}\right]\right)\right) \boldsymbol{d} \phi_{0} \boldsymbol{d} \phi_{i}
\end{aligned}
$$

In this way, we can visualize $P_{i} \cap P_{j}$ as the surface integral over the constructed height function $\nabla \times\left(A_{H^{*}}^{[i]}-A_{H^{*}}^{[i][j]}\right)$.

## III. EXAMPLE: GAIT DESIGN FOR A ROBOT WITH A LONG ACTUATED BODY AND INTERMEDIATE-SIZED LEG

In this section we validate our theory by robot experiments, by using a robophysical approach [1, 19]. We first test the whole-body motion resulting from the two common approaches to design quadrupedal body undulation. We next apply our proposed framework to design whole-body motion and visualize the kinematics of whole-body coordination. Finally, We verify that our theory prediction by robot experiments and RFT simulation.

## A. Experimental set-up

We built an open-loop, servo-driven, 3D-printed robot (450 $\mathrm{g}, \sim 40 \mathrm{~cm}$ long) with a long actuated body and intermediatesized legs (Fig 1), and test its performance on granular media


Fig. 4. Height functions (a) The height function presenting the coordination of snake-like body undulation and leg movement. The solid blue curve represents the optimized gait path. (b) The height function presenting the coordination of lizard-like body undulation and leg movement. The solid blue curve represents the optimized gait path. (c1) The height function presenting the coordination of upper back and leg movement. The solid blue curve represents the optimized gait path. (c2) The height function presenting the coordination of lower back and leg movement. The solid blue curve represents the optimized gait path.
( $\sim 1 \mathrm{~mm}$ diameter poppy seeds). Each leg (front right $F R$, front left $F L$, hind right $H R$ and hind left $H L$ ) has two DoFs, with one controlling the planar shoulder joint angle $\beta_{i}$ and the other controlling the binary ground contact $\delta_{i}$. The body consists of three segments, connected by two independentlyactuated joints (upper back joint angle $\alpha_{1}$ and lower back joint angle $\alpha_{2}$ ), and maintains continuous contact with the ground (Fig. 3).

The robot executes a programmed set of movements (following the gait input) to walk on the loosely-packed poppy seeds. Throughout the experiment, four Optitrack Flex13 cameras record the positions of the infrared-reflective markers on the robot. At the end of each experiment, the robot's final position is identified. A 3-axis gantry system moves to the robot, picks it up and places it back to the starting position. An air-fluidized bed erases the footprints and allows the seeds to be reset into a loosely-packed state.

## B. Sub-system motions

We decompose the motion of the robot into three subsystem motions: the leg sub-system motion, the upper back sub-system motion and the lower back sub-system motion.

1) Leg sub-system motion: We assume that the leg subsystem motion follows a trot gait pattern, a typical leg movements given from prior work [9]. The trot is a two-beat gait with diagonally paired leg movement. The fore right leg is always in phase with the hind left leg while fore left leg is always in phase with the hind right leg. In Fig. 3p, we show that shoulder joint angles and the leg contact patterns of all legs are prescribed by a periodic variable $\phi_{1}$.


Fig. 5. Robot experiment and RFT simulation results (a) The primary data of displacement from robot experiments (solid line with error bar) and RFT simulation (solid line) plotted as a function of gait percentage over two gait cycles for snake-like body undulation (blue), optimal body undulation (red) and lizard-like body undulation (yellow). (b) The displacement is plotted as a function of spatial phase. The red curve with error bar shows the robot experiment results. The blue curve shows the RFT simulation results.
2) Upper back sub-system motion: We prescribe the upper back joint angle, $\alpha_{1}$, as a sinusoidal function of the gait phase $\phi_{2}$, as illustrated in Fig. 30. While the upper back motion system itself cannot generate effective locomotion, it can improve the locomotion performance by coordinating with the leg sub-system motion and the lower back sub-system motion.
3) Lower back sub-system motion: Similar to upper back sub-system motion, we prescribe the lower back joint angle, $\alpha_{2}$, as a sinusoidal function of the gait phase $\phi_{3}$, as illustrated in Fig. 3d. While the lower back motion system itself cannot generate effective locomotion, it can improve the locomotion performance by coordinating with the leg sub-system motion and the upper back sub-system motion.

## C. Whole body motion gait design

The body undulation plays an important role in wholebody motion. In some cases where the leg drag is dominating (lizard-like locomotion), the body undulation will form a standing wave to improve the self-propulsion by leg-ground interaction [11, 27]. In other cases where the body drag is dominating (snake-like locomotion), the body undulation will form a traveling serpenoid wave to generate self-propulsion by body-ground interaction [8]. In this section, we will explore
the role of body undulation for a quadrupedal robot with a long actuated body and intermediate-sized leg.

We assume all the mapping functions are linear, i.e., $\phi_{i}=$ $\phi_{j}+\psi_{i j}$, where $\psi_{i j}$ is the phase offset between the sub-system motion $i$ and $j$. We define the spatial phase as the phase shift between the upper back and lower back, i.e., $\psi_{23}$. The spatial phase characterizes the wave pattern of the body undulation.

In our robot experiments, we find that the optimal spatial phase, intermediate between lizard-like and snake-like, maximizes the displacement. The optimized whole-body motion will leverage the self-propulsion of leg-ground contact and the self-propulsion of body-ground contact.

1) Snake-like body undulation: The long actuated body of the robot can generate self-propulsion from body-ground interaction. Hatton et al. [8] studied the motion of threelink swimmer on granular media, where the body undulation follows the traveling serpenoid wave. We call the traveling serpenoid wave body undulation snake-like body undulation. In snake-like body undulation, we couple the upper and lower back sub-system motions such that they form a traveling serpenoid wave propagated from head to tail. Specifically, the phase of lower back sub-system motion is $\pi / 2$ ahead of the phase of upper back sub-system motion, i.e., the spatial phase of snake-like body undulation is $\psi_{23}^{\text {snake }}=\pi / 2$.

Next, we seek to coordinate the traveling serpenoid wave body undulation with the leg movement, which is characterized by $\psi_{12}^{\text {snake }}$. Following the steps in section II-B2, we can calculate the height function on the shape space composed of $\left[\phi_{1}, \phi_{2}\right]^{T}$ and then design the proper whole-body coordination gait. The height function and the optimized whole-body coordination gait for snake-like body undulation is shown in Fig 4 a. To summarize, in snake-like body undulation, we have $\psi_{12}^{\text {snake }}=0.75 \pi, \psi_{13}^{\text {snake }}=1.25 \pi$ and $\psi_{23}^{\text {snake }}=0.5 \pi$.
2) Lizard-like body undulation: The long actuated body can, on the other hand, help to improve the self-propulsion by leg-ground interaction. Zhong et al. [27] studied the motion of a quadrupedal robot with a single DoF in the back. Ijspeert et al. [11] showed that standing wave body undulation can be coordinated with the leg movement during walking. We call the standing wave body undulation lizard-like body undulation. Although the lizard-like body undulation itself does not lead to effective motion, it can be coordinated with leg movement to improve locomotion performance.

In lizard-like body undulation, we couple the upper and lower back such that they form a standing wave. Specifically, lower back sub-system motion is always in phase with upper back sub-system motion, i.e., $\phi_{2}=\phi_{3}$, such that spatial phase $\psi_{23}^{\text {lizard }}=0$.

Next, we seek to coordinate the standing wave body undulation with the leg movement, which is characterized by $\psi_{12}^{l i z a r d}$. Following the steps in section II-B2, we can calculate the height function on the two-dimensional shape space composed of $\left[\phi_{1}, \phi_{2}\right]^{T}$ and then design the proper wholebody coordination gait. The height function and the optimized whole-body coordination gait for lizard-like body undulation is shown in Fig 4b. To summarize, in lizard-like body undulation,
we have $\psi_{12}^{\text {lizard }}=\pi, \psi_{13}^{\text {lizard }}=\pi$ and $\psi_{23}^{\text {lizard }}=0$.
3) Optimal body undulation: In this subsection, we consider the upper back and lower back as two independent subsystem motions. With the methods described in section II-B2. we can design the whole-body motion gait in the threedimensional sub-system motion phase space. Fig. 4 chows the converged optimal gaits on their corresponding height functions.

In optimal body undulation, we have $\psi_{12}^{\text {optimal }}=1.13 \pi$, $\psi_{13}^{\text {optimal }}=0.88 \pi$ and $\psi_{23}^{\text {optimal }}=0.25 \pi$. We observed that the spatial phase in optimal undulation $\psi_{23}^{\text {optimal }}$ is intermediate between $\psi_{23}^{\text {snake }}$ and $\psi_{23}^{l i z a r d}$.

The theory-predicted optimal body undulation for our robot is neither traveling serpenoid wave nor standing wave. Instead, the intermediate spatial space in the spectrum, from lizardlike to snake-like, will maximize the forward displacement. To verify our theory, we sampled the spectrum and tested the the resulting whole-body motion on the robot experiments and RFT simulation. Both robot experiments and RFT simulation agree well with our theory prediction.

## D. Visualization

In order to visualize the contribution from each individual sub-system motion, we present a Venn diagram for sub-system motions in Table. II There are three major sets: $P_{l e g}$, contribution from the leg movement (blue circle), $P_{u p}$, contribution from the upper back (red circle), and $P_{\text {low }}$, the contribution from the lower back (yellow circle). The intersection of three sets create 4 additional regions (for a total of 7). The physical meaning of each region is listed in Table. []

Note that region 1, 3 and 6 are the contribution solely from upper back, lower back and leg sub-system motion respectively. The contribution solely from an individual sub-system motion is a pre-computed constant, independent from any gait we prescribe. Other regions, presenting the contribution from coordination, are determined by the gait paths (characterize by the phase offset $\psi_{12}, \psi_{13}$ and $\psi_{23}$ ) we choose. In Table. I we listed the dependence of each region on the phase offsets.

Following the steps in Section II-B3, we can visualize the contribution from each region in Fig. 6 Note that in a motion system with three sub-system motions, all the regions in Venn Diagram can be visualized from the basic set operations. From Fig. 6f and Fig. 6d, we observe that optimized body undulation indeed lead to additional displacement by increasing the selfpropulsion from the leg-ground interaction, similar to the role of lizard body undulation. On the other hand, Fig. 6a suggests that the slight phase offset between upper back and lower back, $\psi_{23}$ can lead to additional displacement by self-propulsion from body-leg interaction, similar to the role of snake body undulation.

## E. Results

In our robot experiments, we recorded the displacement of the robot as it implemented the different gaits. In Fig. 5a, we plotted the displacement as a function of gait percentage for the snake-like, optimal body and lizard-like body undulation

| Venn Diagram | Region | Definition | Physical meaning | Determined by |
| :---: | :---: | :---: | :---: | :---: |
|  | 1 | $P_{u p} \cap\left(P_{\text {leg }} \cup P_{\text {low }}\right)^{C}$ | solely from UB | - |
|  | 2 | $\left(P_{\text {up }} \cap P_{\text {low }}\right) \cap P_{\text {leg }}^{C}$ | from coordination of UB and LB | $\psi_{23}$ |
|  | 3 | $P_{l o w} \cap\left(P_{l e g} \cup P_{u p}\right)^{C}$ | solely from LB | - |
|  | 4 | $P_{\text {low }} \cap P_{\text {leg }} \cap P_{\text {up }}$ | whole body coordination | $\psi_{12}$ and $\psi_{13}$ |
|  | 5 | $\left(P_{\text {leg }} \cap P_{\text {low }}\right) \cap P_{u p}^{C}$ | from coordination of LB and legs | $\psi_{13}$ |
|  | 6 | $P_{l e g} \cap\left(P_{u p} \cup P_{\text {low }}\right)^{C}$ | solely from legs | - |
|  | 7 | $\left(P_{\text {leg }} \cap P_{\text {low }}\right) \cap P_{u p}^{C}$ | from coordination of UB and legs | $\psi_{12}$ |

TABLE I
DEFINITION AND PHYSICAL MEANING OF EACH OF THE REGIONS IN THE VENN DIAGRAM FOR OUR ROBOT.
 can be visualized by the surface integral over the corresponding height function.
gaits. In Fig. 5b, we plotted the displacement as a function of spatial phase. Both simulation and robot experiments suggest that our theoretically calculated optimal spatial phase will lead to largest displacement.

We observe that the optimal body undulation spatial phase decreases from $\pi / 2$ to 0 for a robot with a long actuated body and infinitesimal-sized leg (like a snake), a robot with a long actuated body and intermediate-sized legs, and a robot with a long actuated body and long legs (like a lizard). Such spatial phase decreases could be associated with the body size changes. We hypothesize that such changes are associated with the evolution. This framework may provide insight into the hypothesis on the correspondence between form and function [18].

## IV. CONCLUSION

In this paper, we presented a new framework for designing locomotion by coordinating motions of different subsystems, as well as visualizing the kinematics of the resulting whole-body motion. Our framework allows us to design, optimize, and visualize the whole-body motion. We applied our framework to design a whole-body motion of quadrupedal robots. A quadrupedal robot with a long actuated body and infinitesimally small legs (snake-like) mainly utilizes its bodyground interaction to generate self-propulsion. On the contrary, the quadrupedal robot with a long actuated body and long
legs (lizard-like) mainly unitizes its leg-ground interaction to generate self-propulsion. Our quadrupedal robot, with a long actuated body and intermediate-sized legs, combines both approaches to coordinate the self-propulsion from both body-leg interaction and leg-ground interactions to generate efficient locomotion. Our theory-predicted optimized wholebody motion agrees well with the presented robot experiments and RFT simulation.

Although this paper only considered cases where the hierarchy had a depth of one (i.e., no sub-sub-systems), we believe that our approach can be used recursively to optimize locomotion for a robot decomposed into a deeper hierarchy of sub-systems, sub-sub-systems, etc. Additionally, this framework can also be applied to complex high-dimensional motion systems where many sub-systems exist. In future work, we will extend our work to study real-time motion planning in response to possible sub-system failures. That is, when some joints fail to be actuated, we envision that our framework can design a new whole-body motion in response to these new constraints. Additionally, we will investigate how our framework can be used to analyze animal locomotion with numerous sub-systems in complex environments.

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